

Assignment 8

Rec'd Problem 5.1

Adiabatic expansion of ideal gas:

$$(V_i, T_i) \rightarrow (V_f, T_f)$$

a) We can calculate T_f using

$$pV^\gamma = \text{constant}$$

along with the EOS

$$p = \frac{NkT}{V}$$

$$\Rightarrow \frac{NkT}{V} V^\gamma = \text{const.}$$

If $N = \text{const.}$, then

$$\boxed{T V^{\gamma-1} = \text{const.}}$$

Hence we find that

$$T_i V_i^{\gamma-1} = T_f V_f^{\gamma-1}$$

$$\Rightarrow \boxed{T_f = T_i \left(\frac{V_i}{V_f} \right)^{\gamma-1}}$$

b) next we see if we can get the same result starting with the assumption of constant entropy.

$$ds = \left(\frac{\partial s}{\partial T}\right)_V dT + \left(\frac{\partial s}{\partial V}\right)_T dV = 0$$

$$\Rightarrow \left(\frac{\partial T}{\partial V}\right)_s = - \frac{\left(\frac{\partial s}{\partial V}\right)_T}{\left(\frac{\partial s}{\partial T}\right)_V}$$

next we use the second law to write

$$dQ = T ds$$

$$\Rightarrow \left(\frac{\partial s}{\partial T}\right)_V = \frac{1}{T} \left(\frac{\partial Q}{\partial T}\right)_V = \frac{C_V}{T}$$

where C_V is the heat capacity. For an ideal gas,

$$C_V = \left(\frac{\partial E}{\partial T}\right)_V = \left(\frac{\partial \frac{3}{2} NkT}{\partial T}\right)_V$$

$$\Rightarrow \boxed{C_V = \frac{3}{2} Nk}$$

Next we use one of Maxwell's relations to write

$$\left(\frac{\partial s}{\partial V}\right)_T = \left(\frac{\partial P}{\partial T}\right)_V$$

For an ideal gas this yields

$$\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial \frac{NkT}{V}}{\partial T}\right)_V = \frac{Nk}{V}$$

Combining results, we find that

$$\left(\frac{\partial T}{\partial V}\right)_S = - \frac{\frac{Nk}{V}}{\frac{1}{T} \frac{3}{2} Nk} = - \frac{2}{3} \frac{T}{V}$$

$$a \quad \left(\frac{\partial \ln T}{\partial \ln V}\right)_S = - \frac{2}{3}$$

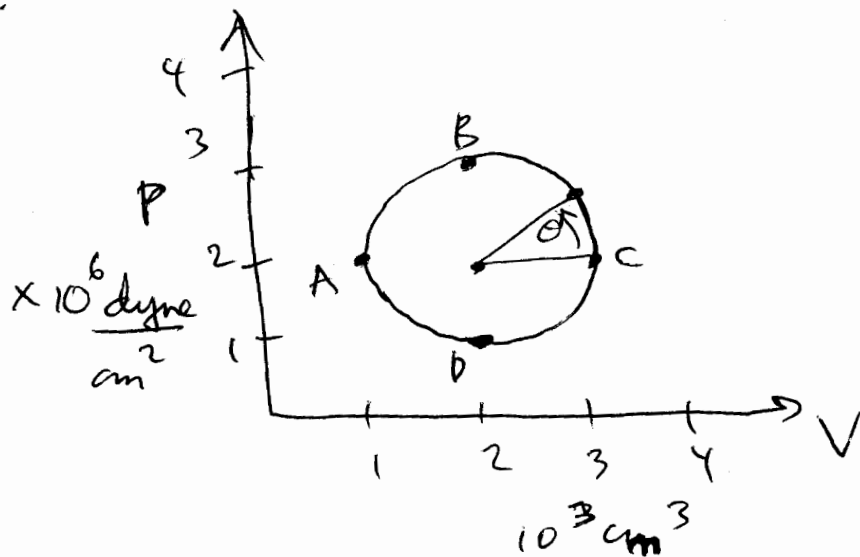
$$\Rightarrow T \propto V^{-2/3}$$

$$\Rightarrow \boxed{TV^{2/3} = \text{constant}}$$

This is consistent with part (a), since $\gamma - 1 = \frac{5}{3} - 1 = \frac{2}{3}$ for an ideal gas.

Problem 5.2

Consider the process indicated in the diagram:



a) work done in one cycle is computed by writing

$$dW = P dV$$

$$\therefore \Delta W = \int P dV$$

We need to transform to polar coordinates using

$$V = (2 + \cos \theta) \text{ cm}^3 \cdot 10^3$$

$$P = (2 + \sin \theta) \frac{\text{dyne}}{\text{cm}^2} \cdot 10^6$$

b) Compute the energy difference between state C and state A. Since this is an ideal gas, we have

$$E = \frac{3}{2} NLT, \quad PV = NLT$$

$$\therefore E = \frac{3}{2} PV$$

Therefore we find that

$$\begin{aligned} E_A &= \frac{3}{2} \cdot 2 \times 10^6 \frac{\text{dynes}}{\text{cm}^2} \cdot 10^3 \text{cm}^3 \\ &= 3 \times 10^9 \cdot 10^{-5} \text{N} \cdot 10^{-2} \text{m} \end{aligned}$$

$$\therefore E_A = 300 \text{ J}$$

Likewise, we find that

$$E_C = \frac{3}{2} \cdot 3 \times 10^6 \frac{\text{dynes}}{\text{cm}^2} \cdot 3 \times 10^3 \text{cm}^3$$

$$\therefore E_C = 9 \times 10^9 \cdot 10^{-5} \text{N} \cdot 10^{-2} \text{m}$$

$$\therefore E_C = 900 \text{ J}$$

The energy difference is therefore

$$E_C - E_A = 600 \text{ J}$$

We obtain

$$\frac{dV}{d\theta} = -\sin\theta \text{ cm}^3 10^3$$

The integral now becomes for one cycle

$$\Delta W = - \int_{2\pi}^0 (2 + \sin\theta) 10^6 \frac{\text{dyne}}{\text{cm}^2} \sin\theta 10^3 \text{ cm}^3 d\theta$$

Note that

$$1 \text{ dyne} = 1 \frac{\text{g cm}}{\text{s}^2} = \frac{10^{-3} \text{ kg } 10^{-2} \text{ m}}{\text{s}^2}$$

$$\therefore 1 \text{ dyne} = 10^{-5} \text{ N} \quad (\text{N} = \text{Newton})$$

We find that

$$\Delta W = 10^9 10^{-5} \text{ N cm} \int_0^{2\pi} (2 + \sin\theta) \sin\theta d\theta$$

$$= 10^4 \text{ N } 10^{-2} \text{ m} \cdot \pi \quad (1 \text{ N} \cdot \text{m} = 1 \text{ J} \leftarrow \text{Joule})$$

\therefore

$$\Delta W = 314 \text{ J}$$

This is the work done by the system in one cycle.

c) The heat absorbed in going from A to C along the path ABC is given by

$$\Delta E = \Delta Q - \Delta W$$

$$\begin{aligned} \Delta W &= -10^4 \text{ N } 10^{-2} \text{ m } \int_{\pi}^0 (2 + \sin \theta) \sin \theta \, d\theta \\ &= 100 \text{ J} \cdot \int_0^{\pi} (2 + \sin \theta) \sin \theta \, d\theta \\ &= 100 \text{ J} \cdot \left(4 + \frac{\pi}{2}\right) \end{aligned}$$

$$\Delta W = 557.08 \text{ J}$$

Then we have

$$\begin{aligned} \Delta Q &= \Delta E + \Delta W \\ &= E_C - E_A + 557.08 \text{ J} \\ &= 600 \text{ J} + 557.08 \text{ J} \end{aligned}$$

$$\Delta Q = 1157 \text{ J}$$

Reif Problem 5.7

The forces of gravity and pressure must balance;



$dz \left\{ \begin{array}{l} A \\ \leftarrow dm = A \rho dz \end{array} \right.$

$\rho =$ mass density,

$$\rho = \frac{N h T}{V}$$

$$= \frac{V R T}{V} \frac{\mu}{\mu}$$

$$= \frac{\rho R T}{\mu}, \quad \rho = \frac{\nu \mu}{V}$$

a) Balance of forces requires that

$$g \, dm = -dp \cdot A$$

Hence

$$\begin{aligned} dp &= - \frac{g A \rho dz}{A} \\ &= -g dz \frac{\rho P}{RT} \end{aligned}$$

$$\therefore \boxed{\frac{dp}{P} = - \frac{g}{RT} dz}$$

b) if the pressure change with height is adiabatic, then from problem 5.5 we have

$$T P^{\frac{1-\gamma}{\gamma}} = \text{const.}$$

$$\Rightarrow T^{\frac{\gamma}{1-\gamma}} P = \text{const.}$$

$$\Rightarrow P \propto T^{\frac{\gamma}{\gamma-1}}$$

$$\Rightarrow \boxed{\frac{dp}{P} = \frac{\gamma}{\gamma-1} \frac{dT}{T}}$$

c) Combining (a) and (c) we find that

$$\frac{\gamma}{\gamma-1} \frac{dT}{T} = - \frac{\mu g}{RT} dz$$

$$\Rightarrow \boxed{\frac{dT}{dz} = - \left(\frac{\gamma-1}{\gamma}\right) \frac{\mu g}{R}}$$

d) If the atmosphere is isothermal, then (a) yields

$$\int_{p_0}^p d \ln p = - \frac{\mu g}{RT} \int_{z_0}^z dz \quad z_0 = \text{height at sea level}$$

$$\Rightarrow \ln(p/p_0) = - \frac{\mu g}{RT} (z - z_0)$$

$$\therefore \boxed{p = p_0 e^{-\frac{\mu g}{RT} (z - z_0)}}$$

e) If the atmosphere is adiabatic, then

$$\frac{T}{T_0} = \left(\frac{p}{p_0}\right)^{\frac{\gamma-1}{\gamma}}$$

Substitution into the result from (a) yields

$$\begin{aligned} \frac{dp}{dz} &= -\frac{\mu g}{R} \frac{p}{T} = -\frac{\mu g}{R} \frac{p}{T_0} \left(\frac{p}{p_0}\right)^{\frac{1-\gamma}{\gamma}} \\ &= -\frac{\mu g}{R} \frac{p_0^{\frac{\gamma-1}{\gamma}}}{T_0} p^{\frac{1}{\gamma}} \end{aligned}$$

$$\Rightarrow \int_{p_0}^p p^{-\frac{1}{\gamma}} dp = -\frac{\mu g}{R} \frac{p_0^{\frac{\gamma-1}{\gamma}}}{T_0} \int_{z_0}^z dz$$

$$\frac{p^{1-\frac{1}{\gamma}}}{1-\frac{1}{\gamma}} \Big|_{p_0}^p = -\frac{\mu g}{R} \frac{p_0^{\frac{\gamma-1}{\gamma}}}{T_0} (z-z_0)$$

$$p^{\frac{\gamma-1}{\gamma}} - p_0^{\frac{\gamma-1}{\gamma}} = -\left(\frac{\gamma-1}{\gamma}\right) \frac{\mu g}{R} \frac{p_0^{\frac{\gamma-1}{\gamma}}}{T_0} (z-z_0)$$

a

$$\boxed{\left(\frac{p}{p_0}\right)^{\frac{\gamma-1}{\gamma}} - 1 = -\left(\frac{\gamma-1}{\gamma}\right) \frac{\mu g}{RT_0} (z-z_0)}$$