

Reif Problem 5.19

The van der Waals equation of state is

$$\left(p + \frac{a}{v^2}\right)(v-b) = RT$$

where $v \equiv \frac{V}{\nu}$ is the molar volume and ν is the number of moles. Recall that the gas constant R is related to Boltzmann's constant k via $\nu R = N k$ and therefore $R = N_A k$ where $N_A =$ Avogadro's number. The critical point is defined by the conditions

$$\left(\frac{\partial p}{\partial v}\right)_T = 0, \quad \left(\frac{\partial^2 p}{\partial v^2}\right)_T = 0$$

We have

$$p = \frac{RT}{v-b} - \frac{a}{v^2}$$

$$\therefore \frac{\partial p}{\partial v} = -\frac{RT}{(v-b)^2} + \frac{2a}{v^3}$$

$$\frac{\partial^2 p}{\partial v^2} = \frac{2RT}{(v-b)^3} - \frac{6a}{v^4}$$

Setting these derivatives equal to zero yields

$$\frac{RT}{(v-b)^2} = \frac{2a}{v^3}, \quad \frac{RT}{(v-b)^3} = \frac{3a}{v^4}$$

Dividing the first by the second gives

$$v-b = \frac{2}{3}v \quad \Rightarrow \quad \boxed{v_c = 3b}$$

We also obtain

$$RT_c = (v-b)^2 \frac{2a}{v^3} = (2b)^2 \frac{2a}{v^3}$$

$$\therefore RT_c = \frac{8b^2 a}{27b^3} = \frac{8a}{27b}$$

$$\Rightarrow \quad \boxed{T_c = \frac{8a}{27bR}}$$

For the pressure we obtain

$$p = \frac{RT}{v-b} - \frac{a}{v^2}$$

$$\therefore p_c = \frac{8a}{27b} \frac{1}{2b} - \frac{a}{9b^2}$$

$$\text{or } p_c = \frac{4a - 3a}{27b^2} \quad \Rightarrow \quad \boxed{p_c = \frac{a}{27b^2}}$$

Combining these results we find that

$$(a) \quad \boxed{\begin{aligned} b &= \frac{v_c}{3} \\ a &= \frac{9v_c T_c R}{8} \end{aligned}}$$

(b) From ~~the~~ equation of state we find that

$$p_c = \frac{a}{27b^2} = \frac{9v_c T_c R}{8} \frac{9}{27v_c^2}$$

$$= \boxed{p_c = \frac{3RT_c}{8v_c}}$$

(c) We can write ~~the~~ equation of state as

$$\left(p' p_c + \frac{a}{(v'v_c)^2} \right) (v'v_c - b) = RT' T_c$$

or

$$\left(p' \frac{3RT_c}{8v_c} + \frac{9v_c T_c R}{8v'^2 v_c^2} \right) (v'v_c - \frac{v_c}{3}) = RT' T_c$$

This reduces to

$$\left(\frac{3}{8} p' + \frac{9}{8v'^2} \right) (v' - \frac{1}{3}) = T'$$

or

$$\boxed{\left(p' + \frac{3}{v'^2} \right) (3v' - 1) = 8T'}$$

Assuming that $C_V = \text{constant}$, we can carry out the integration to obtain

$$T_2 - T_1 = \frac{a v^2}{C_V v} \Big|_{v_1}^{v_2}$$

or

$$T_2 - T_1 = \frac{a v^2}{C_V} \left(\frac{1}{v_2} - \frac{1}{v_1} \right)$$

Assignment 11

Reif Problem 6.1

Oscillator in contact with heat reservoir. This system is governed by the canonical distribution:

$$E_n = \left(n + \frac{1}{2}\right) \hbar \omega$$

$$P_n = C e^{-E_n/\beta} = C e^{-E_n/kT}$$

$$a) \frac{P_1}{P_0} = \frac{C e^{-E_1/kT}}{C e^{-E_0/kT}} = \frac{e^{-\frac{3}{2} \frac{\hbar \omega}{kT}}}{e^{-\frac{1}{2} \frac{\hbar \omega}{kT}}}$$

$$\Rightarrow \frac{P_1}{P_0} = e^{-\hbar \omega/kT}$$

b) If only the ground state and the first excited state are occupied, then the mean energy is given by

$$\bar{E} = P_1 E_1 + P_0 E_0$$

$$= \frac{E_1 e^{-E_1/kT} + E_0 e^{-E_0/kT}}{e^{-E_1/kT} + e^{-E_0/kT}}$$

$$= \frac{\frac{3}{2} \hbar \omega e^{-\frac{3}{2} \frac{\hbar \omega}{kT}} + \frac{1}{2} \hbar \omega e^{-\frac{1}{2} \frac{\hbar \omega}{kT}}}{e^{-\frac{3}{2} \frac{\hbar \omega}{kT}} + e^{-\frac{1}{2} \frac{\hbar \omega}{kT}}}$$

\therefore

$$\bar{E} = \frac{\hbar \omega}{2} \frac{3 e^{-\frac{\hbar \omega}{kT}} + 1}{e^{-\frac{\hbar \omega}{kT}} + 1}$$